Optimal inductor design for nanofluid heating characterisation

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Abstract

Purpose – Magnetic fluid hyperthermia experiment requires a uniform magnetic field in order to control the heating rate of a magnetic nanoparticle fluid for laboratory tests. The automated optimal design of a real-life device able to generate a uniform magnetic field suitable to heat cells in a Petri dish is presented. The paper aim to discuss these issues.

Design/methodology/approach – The inductor for tests has been designed using finite element analysis and evolutionary computing coupled to design of experiments technique in order to take into account sensitivity of solutions.

Findings – The geometry of the inductor has been designed and a laboratory prototype has been built. Results of preliminary tests, using a previously synthesized and characterized magneto fluid, are presented. **Originality/value** – Design of experiment approach combined with evolutionary computing has been used to compute the solution sensitivity and approximate a 3D Pareto front. The designed inductor has been tested in an experimental set-up.

Keywords Multi-objective optimization, Finite element analysis, Sensitivity analysis,

Magnetic fluid hyperthermia, Magnetic nanoparticles, Uniform magnetic field

Paper type Research paper

1. Introduction

Magnetic fluid hyperthermia (MFH) is a cancer therapy that uses a magnetic nanoparticle (NP) fluid to heat locally tumor cells (Hildebrandt *et al.*, 2002; Goya *et al.*, 2003, 2013; Krishnan, 2010). The first device applied for this therapy is described in Gneveckow *et al.* (2004) and Jordan *et al.* (1999, 2001). At the time being, a large variety of magnetic fluids exists (Roca *et al.*, 2009; Sun *et al.*, 2003). In order to verify the magnetic fluid heating capability in biological systems, the synthesized NPs have to be tested on cell cultures (Jordan, 1999). The geometry of the device to generate the magnetic field must be suitably designed in order to achieve a good uniformity in the region where cells are placed. The design of the magnetic field source can be driven by means of various parametric

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Engineering Computations: International Journal for Computer-Aided Engineering and Software Vol. 32 No. 7, 2015 pp. 1-25 © Emerald Group Publishing Limited 02644401 DOI 10.1108/EC-10-2014-0218 or non-parametric optimization algorithms (Barba et al., 2013; Di Barba et al., 2010, 2012a, b, c, 2013, 2014). In the paper, the optimization of the magnetic field in a device to heat a NP magnetic fluid, previously injected inside cells cultured in a Petri dish, is presented. The inductor design is driven by means of NSGA-DOE algorithm, a multi-objective optimization method NSGA based that incorporates a set of uncertain parameters (Di Barba et al., 2014). In fact, the issue of handling a multi-dimensional solution space in a cost-effective way is a prerequisite for selecting a specific algorithm (Riganti Fulginei and Salvini, 2007). In computational electromagnetism, optimization problems are characterized by one or more objective functions which are known in a numerical form only; usually, their computation rely of 2D or 3D finite-element models of the device or system under study. Therefore, the standard assumption of convexity cannot be invoked by any means and the use of gradient-free algorithms is recommended. Despite the statistical performance of an optimization algorithm across the whole class of solvable problem is the same (Wolpert and Macready, 1997), the evolutionary computing algorithms have proven to be effective in identifying the region of global optimum both for both single-objective and multi-objective problems.

In the paper, a 3D objective space is investigated; in particular, three objective functions are minimized: (f_1) the inhomogeneity of the magnetic field strength in the bottom region of the Petri dish, (f_2) the sensitivity of the magnetic field inhomogeneity and (f_3) either the voltage at inductor ends or maximizing the magnetic field value in the bottom of Petri dish. Finally, the designed device has been realized and tested using NP suspensions produced by the chemical laboratory of the Department of Industrial Engineering of Padova University. Nanofluids have been previously characterized in terms of both physical and chemical properties and also in term of heat generation at a given intensity of time-varying magnetic field.

2. Heating problem

In MFH the heat is generated by the interaction of a magnetic fluid, a suspension of superparamagnetic nanometric particles in a fluid, and a time-varying magnetic field. The magnetic field allows the rotation of both the magnetic moments and entire NPs in the biological medium. The power density (in (W/m^3)) generated in the nanofluid can be evaluated by means of the following relation (Rosensweig, 2002):

$$P = \mu_0 \pi \chi''(H) f H^2 \tag{1}$$

where μ_0 is the magnetic permeability of the vacuum, *f* the magnetic field frequency, χ'' the magnetic susceptibility of the NPs, *H* the root means square amplitude of magnetic field in (Am⁻¹). Moreover, also the magnetic susceptibility χ'' (dimensionless) is a function of the magnetic field *H* and Langevin's parameter, ξ (Rosensweig, 2002):

$$\chi'' = \frac{\omega\tau}{\left[1 + (\omega\tau)^2\right]}\chi_0 \tag{2}$$

where ω is the angular frequency of H in (s⁻¹), τ the relaxation time, in (s), of the NPs due to Neel, τ_N , and Brown, τ_B , relaxation time constants, in (s), such that $\tau^{-1} = \tau_N^{-1} + \tau_B^{-1}$ (Rosensweig, 2002). The Brown relaxation time is correlated to the NP rotation in the suspension media and depends on the medium viscosity, η in (Pa·s):

$$\tau_B = 3\eta V_H (k_B T)^{-1} \text{ and } \tau_N = \tau_0 e^{\frac{K V_m}{k_B T}}$$
(3)

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where k_B is the Boltzman constant $(1.38 \times 10^{-23} \text{ (JK}^{-1}))$, T the temperature in (K), V_m is the volume of the magnetic core with diameter D in (m³), V_H is the hydrodynamic volume considering a cover layer with a thick δ in (m³), $(V_m = 4\pi (D/2)^3/3, V_H = (1 + \delta (D/2)^{-1})^3 V_m)$. Whereas χ_0 (dimensionless) is:

$$\chi_0 = \frac{3\chi_i}{\xi} \cdot \left(\operatorname{coth} \xi - \frac{1}{\xi} \right) \tag{4}$$

with

$$\xi = \frac{\mu_0 M_d H V_m}{k_B T} \text{ and } \chi_i = \frac{\mu_0 \phi M_d^2 V_m}{3kT}$$
(5)

where M_d is the domain magnetization of the magnetic material in (Am⁻¹) and ϕ the fraction (dimensionless) of the magnetic NPs in the suspension (Rosensweig, 2002). In particular, ξ and χ_i are dimensionless quantities. Equation (1) evidences the dependence of the power density P on the square of the magnetic field strength. As a consequence, it is necessary to have a homogeneous magnetic field strength in the volume where there are the NPs. The optimal design of the inductor to generate the magnetic field used in the cell cultures experiments is required in order to apply on the cell surface the same intensity of magnetic field and obtain an uniform heating.

3. Direct problem

The direct problem has been solved by means of 2D models and coupled with an optimization algorithm in order to find the geometry that improves the magnetic field uniformity. Finally, the optimal solution found has been modeled also in a 3D domain to assess the end effects of the inductor on the magnetic field uniformity.

3.1 2D analysis

Figure 1(a) shows the cross section of the 2D axy-symmetric geometry of the device. The Petri dish is placed in a thermally insulated box where a water flow maintains the temperature of the system at 37°C. The magnetic field device is a two-turn inductor with five ferrite blocks placed as in Figure 1(a) in order to concentrate the magnetic flux lines. The ferrite elements have been divided into five blocks for two reasons: first of all, the uniformity of magnetic field in the region of interest can be improved by varying independently the vertical position of each ferrite block in the upper part of the device; moreover, the sizes of blocks that have been considered correspond to commercially available parts. Ferrite elements are made of non-linear material (i.e. initial relative permeability between 800 and 4,300, saturation magnetic flux density between 0.39 and 0.5 T). The magnetic field is solved in time-harmonics using a finite element code (FLUX, n.d.). In particular, the problem is solved in terms of phasor of magnetic vector potential, A, coupled with the electric scalar potential, V. When the Coulomb gauge is applied to the magnetic vector potential, i.e. $\nabla \cdot A = 0$, the following coupled equations are solved (Di Barba *et al.*, 2008, 2012; Binns *et al.*, 1992):

$$\nabla \times \frac{1}{\mu} \nabla \times A + j\omega \mu \frac{1}{\rho} A = -\frac{1}{\rho} \nabla V$$
(6)

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$$\nabla \cdot \frac{1}{\rho} (j\omega\mu A + \nabla V) = 0 \tag{7}$$

with μ is the material permeability, ω field pulsation and ρ material resistivity for copper turns.

The problem has been studied supplying the inductor by imposing either a current or a voltage. A typical mesh (Figure 1(b)) is composed of 23,000 nodes and 13,000 second-order elements.

3.2 3D analysis

Results obtained after the optimization process based on 2D models have been assessed *a posteriori* by means of the 3D model shown in Figure1(c). This model includes the two

inductor turns, represented as two circular bends, and the ferrite elements. The model incorporates also the inductor ends, connected to the power supply. This model disregards Petri and thermal insulating box, because they have no influence on the magnetic field computation. The aim of this model is to assess the influence of the inductor ends on the uniformity of the magnetic field distribution. The problem is solved again in time-harmonics conditions using finite element analysis. The 3D solutions have been obtained by applying the scalar potential formulation in terms of scalar magnetic potential, Φ , and vector electric potential T. The electromagnetic solution is coupled with an external circuit to apply the same current to each turn of the inductor. This formulation allows to compute the magnetic field H and current density J as (Biro *et al.*, 2004):

$$H = T + T_0 - \nabla \Phi \tag{8}$$

$$J = \nabla \times T + \nabla \times T_0 \tag{9}$$

 T_0 is an impressed electric vector potential depending on the imposed current in the coil that is computed as Biot-Savart field.

The ferrites are non-conductive regions with non-linear magnetic properties, where $\nabla \times H = 0$ and the solution is obtained from the total scalar magnetic Φ , by solving:

$$\nabla \cdot (\mu(\boldsymbol{H}) \nabla \Phi) = 0 \tag{10}$$

In the inductor, the reduced scalar magnetic potential is coupled with the electric vector potentials, T and T_0 is calculated by means of Biot-Savart law. The air region that encompasses the active parts in the model is solved in terms of reduced scalar potential. The appropriate coupling conditions between total and reduced scalar potential are applied on the interfaces between air and ferrite (Aliferov *et al.*, 2010; Biro *et al.*, 2004; Meunier, 2008).

The 3D mesh has 5.2×10^5 nodes and 3.0×10^6 first-order volume elements.

4. Inverse problem

The optimization problem aims at minimizing the inhomogeneity of the magnetic field, H, on the bottom of the Petri dish (f_1) , evaluated as in Di Barba *et al.* (2014a, b, in press a, b) with a tolerance interval of ± 10 A/m. The (f_1) is evaluated in a fixed grid of sampling points which does not change during the optimization process and accounts for the number of points that satisfy the "proximity criterion". Specifically, the value of (f_1) is the minimum number of points of the fixed grid for which the magnetic field strength is out of the prescribed tolerance interval. The (f_2) is the sensitivity of (f_1) with respect to a set of uncertain parameters, which are different from the design variables, computed by means of the NSGA-DOE algorithm (Di Barba et al., 2014). The inhomogeneity of H is evaluated on a fixed grid of points positioned in the bottom of the Petri dish; (f_1) is dimensionless and evaluate the maximum number of points out of the tolerance interval. The (f_2) objective function computes the variation of (f_1) due to a variation of the Petri dish position, for a given device geometry (Di Barba et al., 2014a, b, in press a, b). In the NSGA-DOE algorithm the updating of the design variables is managed by means of a NSGA-II inspired algorithm, whereas the uncertain parameters are taken into account by means of a design of experiment (DOE) strategy. Therefore, NSGA-DOE strategy is a combination of NSGA strategy, to sort the current individuals in local fronts, and DOE strategy, to compute the Optimal inductor design

sensitivity (f_2) which in turn is one of the objective function. It is emphasized that uncertain parameters are not coincident with design variables and are used for the computation of the sensitivity of a solution. As far as the third objective function (f_3) two cases have been considered: in the first case, the inductor is supplied by a current of 500 Arms at 350 kHz, and (f_3) is the voltage at the inductor ends to be minimized. In fact, the applied voltage should not exceed the typical value made available at the converter output. In the second case the inductor is supplied by a voltage of $600 V_{rms}$ at 350 kHz and (f_3) is the magnetic field value in the bottom of Petri dish to be maximized (typical values are 6-10 kA/m). In particular, considering the voltage at the inductor ends as the third objective function makes the optimization problem more oriented to real-life engineering; this way, in fact, the operation of the power supply converter is taken into account. On the other hand, after many numerical experiments, it could be stated that the proposed method of multi-objective optimization is able to deal successfully with more than two objective functions. In both cases the same set of eight design variables is searched for. The eight design variables (Figure 1(a)) are the coil distance from the insulated box, "dR", the size of ferrite block under the box, "HF" and "LF", the vertical distance of the four ferrite blocks from the insulated box, "hf₀", "hf₁", "hf2", the vertical size of the turns and their gap, "HS" and "ST". In turn, three uncertain parameters are defined as the Petri dish radial position, "DP", the vertical distance of insulated box from lower ferrite, "SPF", and the vertical position of the Petri dish from the insulated box bottom, "HP". Table I reports the design variables and their variation ranges and the nominal values of parameters and the assigned uncertainty range.

The rows of Table II gives the four Y_j , j = 1, ..., 4, combinations of uncertain parameters to evaluate the sensitivity according to Plackett-Burman DOE strategy (Di Barba *et al.*, 2014; Plackett and Burman, 1946). The "+" and "-" mean the sign alternance of the variation applied to each parameter.

So, given a solution, four $f_{i,j}$, j = 1, ..., 4, $(f_{i,1}, f_{i,2}, f_{i,3}, f_{i,4})$ values of the f_i objective functions are derived by varying the design variable values as in the following: for the *k*th uncertain parameter, k = 1, 3, the sums of f_i values $(f_{i,j})$ corresponding to a "+" in Table I, $S_{+,Pk}$, and the ones corresponding to a "-", $S_{-,Pk}$, are computed as follows:

$$S_{+,Pk} = \sum_{j=1}^{4} f_{i,j} \quad Y_j = '+'$$
(11)

$$S_{-,Pk} = \sum_{j=1}^{4} f_{i,j} \quad Y_j = -'$$
(12)

	Design variable	Range (mm)	Uncertain parameters	Nominal value \pm uncertainty range (mm)
Table I.	dR	(3, 10)	SPF	3±1
Lower and upper	ST	(0, 30)	DP	0 ± 1
bounds of design	HS	(10, 60)	HP	3 ± 1
variables and	LF	(20, 80)		
nominal values and	HF	(5, 20)		
uncertainty intervals	hf0	(1, 30)		
of uncertain	hf1	(1, 30)		
parameters	hf2	(1, 30)		

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Then, the influence s_{Pk} of a variation of the *k*th design variable on the value of objective function f_i is evaluated as [34];

$$s_{Pk} = \frac{S_{+,Pk}}{N_{+}} - \frac{S_{-,Pk}}{N_{-}}$$
(13)

where N_+ and N_- are the number of sign "+" and "-" in the column corresponding to the considered variable in Table I. For the *k*th design variable the partial sensitivity, s_{Pk} , is estimated using (13). Finally, the total sensitivity with respect to all the uncertain design variable, f_2 , is given by:

$$f_2 = \sqrt{\sum_{k=1}^3 s_{Pk}^2}$$
(14)

Under a multi-objective context, f_2 given by (14) can well be regarded as a secondary objective function in addition to the design criterion f_1 .

5. Results

The optimal design results of the magnetic device to heat cell cultures, as well as preliminary experimental results to evaluate the heating of a NP magnetic fluid, are presented.

5.1 Magnetic device design

Figure 2 shows an approximation of the 3D Pareto front obtained for the current-supplied model after running a population of 20 individuals for 50 generations. The initial population was randomly generated. The solutions of different runs have been collected.

Figure 3 reports the 2D orthogonal projections of the Pareto front. In Table III a set of solutions on Pareto front are reported. The corresponding geometries are shown in Figure 4.

Figure 5(a) reports the magnetic field strength on the bottom of the Petri dish as a function of position for the cases in Table III. The device geometry with relevant field lines is shown in Figure 5(b) for the solution S_0 .

Table IV reports the f_1 , named f_1^* , computed with the optimized design variables rounded to the nearest integer. The influence of design variables variation on f_1 is shown in the Table V. The sensitivity of f_1 is calculated by varying only one design variable at time inside the interval $\pm d$, d = 1 mm:

$$\frac{\partial f_1}{\partial v_i} \cong \frac{f_1(x_i+d) - f_1(x_i-d)}{2d} \left[\frac{1}{mm}\right] \tag{15}$$

SPF	DP	HP		
+	+	+	$f_{\rm i,1}$	
—	+	-	$f_{i,2}$	Table II.
—	-	+	$f_{i,3}$ Sign alter	ernance for
+	_	-	$f_{i,4}$ DOE evaluation	lation of f ₂

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Notes: (a) Sensitivity; (b) voltage vs H inhomogeneity

Figure 3. 2D projections of pareto fronts

Figure 6 shows an approximation of the 3D Pareto front obtained for the voltage-supplied model after running a population of 20 individuals for 50 generations. The initial population was random generated. The solutions of different runs have been collected.

Figure 7 reports the 2D orthogonal projections of the Pareto front. In Table V a set of solutions on Pareto front are reported. The corresponding geometries are shown in Figure 8.

Figure 7 (a) reports the magnetic field strength on the bottom of the Petri dish as a function of position for the cases in Table VI. The device geometry with relevant field lines is shown in Figure 9 (b) fort solution S_0 .

Table VII reports f_1^* , i.e. the values of f_1 computed with the optimized design variables rounded to the nearest integer. The influence of design variables variation on f_1 is shown in the Table VIII. The sensitivity of f_1 is calculated using the (15).

The aim of the proposed multi-objective optimization is to give to the designer the possibility to select the best feasible solutions in terms of field uniformity and solution robustness in accordance with the power supply characteristics. The designer can decide to achieve an excellent field uniformity using a weak solution in terms of sensitivity or vice versa accordingly to the experience in practical realizing such a kind of devices. The final choice depends also on the quality of the available manufacturing technology. A picture of the two-turn inductor which has been realized is shown in Figure 10: the internal diameter of 18 cm.

5.2 3D model

The magnetic field has been sampled along some lines on a ring positioned (internal radius 18 mm, external radius 62 mm) in the same position of the Petri dish bottom as in Figure 11.

	dR	st	LF	hf	hf_0	hf_1	hf_2	HS	f_1	f_2	f_3	
S0	3.5	6.5	67.7	7.7	23.7	26.2	22.4	58.8	288	60.1	647.8	
S1	4.7	30.0	43.7	5.6	24.0	26.8	18.9	59.7	819	54.4	588.6	
S2	7.8	8.9	61.0	8.6	24.5	23.8	28.6	38.0	1	108.3	824.1	
S3	5.3	2.5	78.7	13.9	10.2	12.2	11.0	46.5	122	47.1	819.8	
Note	Note: Design variables on $[mm]$, objective function f_1 and f_2 values dimensionless, f_3 in $[V]$											





Table III. Some solutions on the Pareto front

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		dR	st	LF	hf	hf_0	hf_1	hf_2	HS	f_1^*	$f_1 - f_1^*$
	S_0^*	4	7	68	8	24	26	22	59	315	-29
Table IV.	S_1^*	5	30	44	6	24	27	19	60	832	-12
f_1^* obtained rounding	S_2^*	8	9	61	9	25	24	29	38	6	-5
design variables	S_3^*	5	3	79	14	10	12	11	47	154	-33

Figure 12(a) reports the magnetic field strength along the arcs A and B. Any influence of the inductor tails appears evident. Moreover, any influence to the uniformity of the magnetic field strength due to inductor tails appears along lines from C to F (Figure 12(b)).

5.3 Experimental results

Before testing the designed inductor to heat cell cultures previously incubated with NPs, a preliminary experiment has been conducted to evaluate the heat generated in different magnetic fluids named samples A, B and C, respectively (Sun *et al.*, 2003; Sun and Zeng, 2002; Bertani *et al.*, 2015, 2010; Sgarbossa *et al.*, 2011; Scaffaro *et al.*, 2011).

Sample A has been prepared by thermal decomposition of iron(III) chloride (FeCl₃·6H₂O) in 2-pyrrolidone at 270°C for 24 h according to the method proposed by Li *et al.* (2005). Similarly, sample B has been synthesized from a 2:1 mixture of iron(III) acetylacetonate ([Fe(acac)₃]) and manganese(II) acetylacetonate ([Mn(acac)₂]). Sample C has been obtained by thermal decomposition of iron(III) acetylacetonate ([Fe(acac)₃]) in diphenyl ether at 200°C for 2 h and then 300°C for 1 h according to the method developed by Sun *et al.* (2003) for the preparation of NPs with a diameter of 6 nm. The morphology of the NPs has been investigated by transmission electron microscopy (TEM) analysis using a TECNAI FEI G2 microscope. The magnetic properties of the samples have been detected by means of a RSO-SQUID magnetometer. Magnetization (M) vs temperature (T) measurements have been performed in the 6-300 K range using both in ZFC (zero-field-cooled) and FC (field-cooled) configurations

	dR	st	LF	hf	hf_0	hf_1	hf_2	HS
$\begin{array}{c} S_0\\S_1\\S_2\\S_3\end{array}$	-26.5 9.5 14 25	-0.5 -19 0.5 -30.5	33.5 37.5 50 -2.5	$70 \\ -29.5 \\ 14.5 \\ -109$	13 -0.5 -39 -62.5	-76.5 42 6 14.5	-15.5 2.5 -30.5 14	$-6 \\ -8.5 \\ -25 \\ 25.5$



Figure 6. Generated individuals (gray points) and approximated pareto front (black crosses)

Table V. sensitivity of f_1^*

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Notes: (a) Sensitivity; (b) magnetic field strength vs H inhomogeneity



(D'Addato et al., 2010). Magnetization loops were recorded at 300 K. In order to gain indications about the type of magnetic material synthetized, we made resort to Mössbauer spectroscopy (Morais, 2008). Indeed, this technique may probe the chemical, structural and magnetic environment surrounding Fe atoms. Mössbauer data have been collected at room temperature using a ⁵⁷Co in Rh source and the spectrometer has been calibrated using an α -Fe foil.

2D projections of Pareto fronts

Figure 8. Geometries of the designed device on pareto front

In Figure 13, TEM images show that the NPs of sample A and B are flower-shaped; regarding the NPs of sample C, the majority are spherical, but a small fraction of triangle-shaped NPs are found, as well. Both samples A and B exhibit NPs with a radius close to 12 nm, whereas the average radius of NPs of sample C is about 3 nm. NPs sizes have been estimated using ImageJ software ("ImageJ", n.d.) measuring the diameter of a large number of particles in TEM images; standard deviation is intended as computed on the logarithm of the diameter values according to Rosensweig (2002). The Mössbauer data collected on the different samples are presented in Figure 14(a). The spectrum of sample B clearly shows the presence of a magnetic signal, mainly made of two different sextets; the corresponding hyperfine magnetic fields (HF) are $B_{HF,1} = 45.0 \pm 0.1 \text{ T}$ and $B_{HF,2} = 48.63 \pm 0.05 \text{ T}$. As this sample was produced using both Fe and Mn diketonates with a 2:1 ratio, for comparison we may use the HF values measured for magnetite (Diamandescu et al., 1998) and for the natural and synthetic Manganese Ferrite (namely Jacobsite, MnFe₂O₄) (Araujo et al., 1989; Maia et al., 1993). The found HF values and the relative ratio of the sextets areas are closer to those expected for magnetite, indicating that the NPs relative Mn concentration is smaller than that found in the starting mixture.

In samples A and C, no sextet structure is found, but just a doublet structure can be observed; regarding sample C, the doublet structure displays a large width, featuring the effects of possible relaxation phenomena. The hyperfine parameters of sample A and C doublets are in agreement with those found in magnetite, supporting the fact that the NPs are made of magnetite. The absence of sextet structures and the appearance of a doublet structure can be generally traced back to the presence of small size NPs (Spizzo *et al.*, 2004).

Regarding sample C, this observation is in agreement with TEM observations. In the case of sample A, this result may be possibly accounted for considering the presence of the flower-shaped NPs, namely clusters made of smaller NPs. The ZFC-FC measurements of the three samples are presented in Figure 14(b). M vs T data for sample C show a behavior that resembles the superparamagnetic one, as expected for small sized particles. Differently, in samples A and B the signature of magnetic irreversibility extends up to room temperature, as expected for large particles or for systems where the magnetic size of the particles is large due to interparticle interactions (Allia *et al.*, 2006), as it may happen in flower-shaped NPs. From room temperature magnetization measurements we obtained the bulk saturation magnetization, M_d ($M_d = M[emu/g] \cdot d(kg/m^3)$, $1(emu/g) = 1(Am^2/kg)$, assuming for magnetite a bulk density, *d*, of 5,180 kg/m³ and for jacobsite a 4,870 kg/m³ (Rosensweig, 2002; Haynes, 2014); M_d is 109.8 kA/m for the sample A, 85.5 kA/m for sample B and 255.9 kA/m for sample C, whereas the anisotropy constant K is close to 10.45 kJ/m³ for the sample A and 126.5 kJ/m³ for sample C. The strong variation of K for the three

	dR	st	LF	hf	hf_0	hf_1	hf_2	HS	f_1	f_2	f_3	
S_0	5.1 3.0	10.9	79.9 21.7	10.2 15.0	4.6 15.9	26.6 12 9	5.8 5.9	51.1 55.8	72 805	230.4	6,369.4 6 802 7	
S_1 S_2 S_2	3.0 4.7 4.7	23.1 9.9	25.6 78.5	16.4 8.8	1.2 3.4	14.8 26.5	19.1 8.0	53.2 29.9	1,091 359	1.7 8.2	7,246.4 6,451.6	
Note	5_3 4.7 9.9 78.5 8.8 3.4 26.5 8.0 29.9 359 8.2 6,451.6 Note: Design variables on [mm], objective functions f_1 and f_2 dimensionless, f_3 in [A/m]											

Table VI. Some solutions on pareto front

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inductor



samples may be interpreted as a signature of the presence of interparticle magnetic interactions (Allia and Tiberto, 2011). The sample A was suspended in water that has a viscosity close to 0.001 Pa s, whereas the sample C is dispersed in PEG 400 with a viscosity of 90 mPa s. Both sample A and C have a volume fraction ϕ of 0.0019 (10 mg of powder was dispersed in each ml of suspension fluid, $\phi = (mg/ml)/d(kg/m^3)$), whereas for sample B ϕ is 0.0021. The data of the NP samples are summarized in Table IX.

The power density generated by NP suspensions was evaluated before applying them to cell cultures. A sample of 2 ml of a water suspension (NP concentration: 10 mg/ml) placed in a test tube was subjected for 5 min to a magnetic field (maximum intensity 13.6 kA/m at 177 kHz, current of 300 Arms) generated by means of a cylindrical inductor with seven turns, 8 cm internal diameter, connected to an EASYHEAT LI

	dR	st	LF	hf	hf_0	hf_1	hf_2	HS
$egin{array}{c} S_0^* \ S_1^* \ S_2^* \ S_3^* \end{array}$	-24 -5.5 -1.5 -12	-39.5 -4 1 49	$12.5 \\ -7 \\ 0 \\ -1.5$	7.5 2 0 -0.5	64.5 3 -1 -114	48.5 11 0 -91	-27 -16.5 1 70.5	-23.5 1 0 39





for Petri dish experiments

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Table VIII. sensitivity of f_1^*



Figure 11. XY section of the 3D geometry and position of the lines where the magnetic field has been sampled

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Notes: (a) Arc A and B; (b) along lines from C to F (Figure 10)

5,060 converter (Figure 15) or a two turns 18 cm internal diameter (Figure 10). Temperature was measured by means of an Optocom Fotemp-1H thermometer with a TS3/2 fiber optic.

The experimental heating rate has been estimated as the ratio between the temperature increasing, ΔT (K), occurred during the experiment and the time interval of the magnetic field application, Δt (s):

$$HR_m = \frac{\Delta T}{\Delta t} (K/s) \tag{16}$$

The heating rate obtained using (12) has been compared with the expected one, HR_c , estimated in adiabatic conditions according to Fortin *et al.* (2007) and Rosensweig (2002) from the power density *P*, in (W/m³), computed using (1) and taking into account the polydispersion of NPs radii as in Fortin *et al.* (2007) and Rosensweig (2002):

$$HR_c = \frac{P}{d \cdot c} (K/s) \tag{17}$$



Notes: (a) Sample A; (b) sample B; (c) sample C



where c and d are, respectively, the specific heat (J/(kg K)) and the density (kg/m³) of the suspension fluid. In this case the density of the magnetic nanofluid has been approximated with the one of the suspension fluid as in Fortin *et al.* (2007). The measured and computed temperature increasing for the sample described in Table IX are presented in Table X. Figure 16 shows the heating rate (17) for the sample A, B and sample C as a function of the NPs radius considering the nanofluid characteristics of Table IX. The dot represents the computed heating rate for the NPs with the radius in Table IX.

The temperature increasing produced in a 4 ml NP suspensions (sample B) in a Petri dish with a diameter of 3 cm has been measured for 5 min applying a magnetic field (using the inductor in Figure 10 with a maximum magnetic field strength

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of 6.3 kA/m at 240 kHz (supply current 450 Arms) or 7.8 kA/m at 240 kHz (supply current 500 Arms)). In the preliminary experiment with Petri dish no thermal insulation has been used; therefore, the heating rate has not been computed because measurements have not been conducted in adiabatic conditions. In this preliminary experiment 3 ml of nanofluid has been put in the Petri dish.

From reported data it appears that suspensions prepared using synthetized NPs are promising to heat cells.

Figure 14. (a) Mössbauer spectra of samples A (in the inset), B and C; the straight line reported on B sample data is a fit performed according to the model described in the text (b) ZFC and FC magnetization measurements of samples A, B and C



	Sample	R (nm)	σ	M (emu/g)	M _d (kA/m)	K (kJ/m ³)	η (Pas)	φ
Table IX.	A	12.0	0.29	21.2	109.8	10.45	0.001	0.0019
Magnetic nanofluid	B	14.4	0.31	18.0	87.7	7.19	0.001	0.0021
characteristics	C	3.43	0.21	49.4	255.9	126.5	0.09	0.0019

Figure 15. Inductor for preliminary experimental tests



6. Conclusion

Numerical models coupled with optimization algorithms have been applied to the robust design of an electromagnetic device able to produce a uniform magnetic field in a prescribed region. Some preliminary experiments have been also carried out in order to exploit the required performance of the system, in order to properly heat magnetic nanofluids. Finally, the designed device has been realized and tested.

ΔT_{m} (K)
10.8
20
0.1
2.7
3.7
•

Note: Temperature measurements related to sample B_2 and B_3 are performed putting 3 ml of nanofluid (10 mg/ml) in a Petri dish

Table X.Experimental results

Figure 16. Heating rate

computed using (13)

as a function of

nanoparticle radii



Notes: (a) Sample A; (b) sample B; (c) sample C

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