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Modified Structure of Protons and Neutrons in Correlated Pairs

The CLAS Collaboration

The atomic nucleus is made of protons and neutrons (nucleons), that are themselves composed of quarks and gluons. Understanding how the quark-gluon structure of a nucleon bound in an atomic nucleus is modified by the surrounding nucleons is an outstanding challenge. Although evidence for such modification, known as the EMC effect, was first observed over 35 years ago, there is still no generally accepted explanation of its cause [1–3]. Recent observations suggest that the EMC effect is related to close-proximity Short Range Correlated (SRC) nucleon pairs in nuclei [4, 5]. Here we report the first simultaneous, high-precision, measurements of the EMC effect and SRC abundances. We show that the EMC data can be explained by a universal modification of the structure of nucleons in neutron-proton (*np*) SRC pairs and present the first data-driven extraction of this universal modification function. This implies that, in heavier nuclei with many more neutrons than protons, each proton is more likely than each neutron to belong to an SRC pair and hence to have its quark structure distorted.

We study nuclear and nucleon structure by scattering high-energy electrons from nuclear targets. The energy and momentum transferred from the electron to the target determines the space-time resolution of the reaction, and thereby, which objects are probed (i.e., quarks or nucleons). To study the structure of nuclei in terms of individual nucleons, we scatter electrons in quasi-elastic (QE) kinematics where the transferred momentum typically ranges from 1 to 2 GeV/c and the transferred energy is consistent with elastic scattering from a moving nucleon. To study the structure of nucleons in terms of quarks and gluons, we use Deep Inelastic Scattering (DIS) kinematics with larger transferred energies and momenta.

Atomic nuclei are broadly described by the nuclear shell model, in which protons and neutrons move in well-defined quantum orbitals, under the influence of an average mean-field created by their mutual interactions. The internal quark-gluon substructure of nucleons was originally expected to be independent of the nuclear environment because quark interactions occur at shorter-distance and higher-energy scales than nuclear interactions. However, DIS measurements indicate that quark momentum distributions in nucleons are modified when nucleons are bound in atomic nuclei [1, 2, 6, 7], breaking down the scale separation between nucleon structure and nuclear structure.

This scale separation breakdown in nuclei was first observed thirty-five years ago in DIS measurements performed by the European Muon Collaboration (EMC) at CERN [8]. These showed a decrease of the DIS cross-section ratio of iron to deuterium in a kinematical region corresponding to moderate- to high-momentum quarks in the bound nucleons. The EMC effect has been confirmed by subsequent measurements on a wide variety of nuclei, using both muons and electrons [9, 10], and over a large range of transferred momenta, see reviews in [1, 2, 6, 7]. The maximum reduction in the DIS cross-section ratio of a nucleus relative to deuterium increases from about 10% for ^{4}He to about 20% for Au.

The EMC effect is now largely accepted as evidence that quark momentum distributions are different in bound nucleons relative to free nucleons [1, 2, 7]. However, there is still no consensus as to the underlying nuclear dynamics driving it.

Currently, there are two leading approaches for describing the EMC effect, which are both consistent with data: (A) all nucleons are slightly modified when bound in nuclei, or (B) nucleons are unmodified most of the time, but are modified significantly when they fluctuate into SRC pairs. See Ref. [1] for a recent review.

SRC pairs are temporal fluctuations of two strongly-interacting nucleons in close proximity, see e.g. [1, 11]. Electron scattering experiments in QE kinematics have shown that SRC pairing shifts nucleons from low-momentum nuclear shell-model states to high-momentum states with momenta greater than the nuclear Fermi momentum. This "high-momentum tail" has a similar shape for all nuclei. The relative abundance of SRC pairs in a nucleus relative to deuterium approximately equals the ratio of their inclusive (e,e′) electron scattering cross-sections in selected QE kinematics [12–15].

Recent studies of nuclei from ^{4}He to Pb [16–22], showed that SRC nucleons are "isophobic"; i.e., similar nucleons are much less likely to pair than dissimilar nucleons, leading to many more *np* SRC pairs than neutron-neutron (*nn*) and proton-proton (*pp*) pairs. The probability for a neutron to be part of an *np*-SRC pair is observed to be approximately constant for all nuclei, while that for a proton increases approximately as *N*/*Z*, the relative number of neutrons to protons [22].

The first experimental evidence supporting the SRC-modification hypothesis as an explanation for the EMC effect came from comparing the abundances of SRC pairs in different nuclei with the size of the EMC effect. Not only do both increase from light to heavy nuclei, but there is a robust linear correlation between them [4, 5]. This suggests that the EMC effect might be related to the high-momentum nucleons in nuclei.

The analysis reported here was motivated by the quest to understand the underlying patterns of nucleon structure modification in nuclei and how this varies from symmetric to asymmetric nuclei. We measured both the DIS and QE inclusive cross-sections simultaneously for deuterium and heavier nuclei, thereby reducing the uncertainties in the extraction of the EMC effect and SRC scaling factors. We observed that: (1) the EMC effect in all measured nuclei is consistent with being due to the universal modification of the internal structure of nucleons in *np*-SRC pairs, permitting the first data-driven extraction of this universal modification function, (2) the measured per-proton EMC effect and SRC probabilities continue to increase with atomic mass A for all measured nuclei while the per-neutron ones stop increasing at $A \approx 12$, and (3) the EMC-SRC correlation is no longer linear when the EMC data are not corrected for unequal numbers of proton and neutrons. We also constrained the internal structure of the free neutron using the extracted universal modification function and we concluded that in neutron-rich nuclei the average proton structure modification will be larger than that of the average neutron.

We analyzed experimental data taken using the CLAS spectrometer [23] at the Thomas Jefferson National Accelerator Facility (Jefferson Lab). In our experiment, a 5.01 GeV electron beam impinged upon a dual target system with a liquid deuterium target cell followed by a foil of either C, Al, Fe or Pb [24]. The scattered electrons were detected in CLAS over a wide range of angles and energies which allowed extracting both QE and DIS reaction cross-section ratios over a wide kinematical region (See Supplementary Information section I).

The electron scattered from the target by exchanging a single virtual photon with momentum $\vec{q}$ and energy ν, giving a four-momentum transfer $Q^2 = |\vec{q}|^2 - \nu^2$. We used these variables to calculate the invariant mass of the nucleon plus virtual photon $W^2 = (m+\nu)^2 - |\vec{q}|^2$ (where m is the nucleon mass) and the scaling variable $x_B = Q^2/2m\nu$.

We extracted cross-section ratios from the measured event yields by correcting for experimental conditions, acceptance and momentum reconstruction effects, reaction effects, and bin-centering effects. See Supplementary Information section I. This was the first precision measurement of inclusive QE scattering for SRCs in both Al and Pb, as well as the first measurement of the EMC effect on Pb. For other measured nuclei our data are consistent with previous measurements but with reduced uncertainties.

The DIS cross-section on a nucleon can be expressed as a function of a single structure function, $F_2(x_B, Q^2)$. In the parton model, x_B represents the fraction of the nucleon momentum carried by the struck quark. $F_2(x_B, Q^2)$ describes the momentum distribution of the quarks in the nucleon, and the ratio, $[F_2^A(x_B, Q^2)/A] \,/\, [F_2^d(x_B, Q^2)/2]$, describes the relative quark momentum distributions in nucleus A and deuterium [2, 7]. For brevity, we will often omit explicit reference to x_B and Q^2, i.e., writing F_2^A/F_2^d, with the understanding that the structure functions are being compared at identical x_B and Q^2. Because the DIS cross-section is proportional to F_2, experimentally the cross-section ratio of two nuclei is assumed to equal their structure-function ratio [1, 2, 6, 7]. The magnitude of the EMC effect is defined by the slope of either the cross-section or the structure-function ratios for $0.3 \le x_B \le 0.7$ (see Supplementary Information sections IV and V).

Similarly, the relative probability for a nucleon to belong to an SRC pair is interpreted as equal to a_2, the average value of the inclusive QE electron-scattering per-nucleon cross-section ratios of nucleus A compared to deuterium at momentum transfer $Q^2 > 1.5$ GeV^2 and $1.45 \le x_B \le 1.9$ [1, 11-15] (see Supplementary Information section III).

Other nuclear effects are expected to be negligible. The contribution of three-nucleon SRCs should be an order of magnitude smaller than the SRC pair contributions. The contributions of two-body currents (called "higher-twist effects" in DIS scattering) should also be small (see Supplementary Information section VIII).

Figure 1 shows the DIS and QE cross-section ratios for scattering off the solid target relative to deuterium as a function of x_B. The red lines are fits to the data that are used to determine the EMC effect slopes or SRC scaling coefficients (see Extended Data Table I and II). Typical 1σ cross-section ratio normalization uncertainties of 1 – 2% directly contribute to the uncertainty in the SRC scaling coefficients but introduce a negligible EMC slope uncertainty.

None of the ratios presented have isoscalar corrections (cross-section corrections for unequal numbers of protons and neutrons), in contrast to much published data. We do this for two reasons, (1) to focus on asymmetric nuclei and (2) because the isoscalar corrections are model-dependent and differ among experiments [9, 10] (see Extended Data Fig. 1).

The DIS data was cut on $Q^2 > 1.5$ GeV^2 and $W > 1.8$ GeV, which is just above the resonance region [25] and higher than the $W > 1.4$ GeV cut used in previous JLab measurements [10]. The extracted EMC slopes are insensitive to variations in these cuts over Q^2 and W ranges of $1.5 - 2.5$ GeV^2 and $1.8 - 2$ GeV respectively (see Supplementary Information Table VII).

Motivated by the correlation between the size of the EMC effect and the SRC pair density (a_2), we model the modification of the nuclear structure function, F_2^A, as due entirely to the modification of *np*-SRC pairs. F_2^A is therefore decomposed into contributions from unmodified mean-field protons and neutrons (the first and second terms in Eq. 1), and *np*-SRC pairs with modified structure functions (third term):

$$\begin{aligned} F_2^A &= (Z - n_{SRC}^A)F_2^p + (N - n_{SRC}^A)F_2^n + n_{SRC}^A(F_2^{p*} + F_2^{n*}) \qquad \text{Eq. 1} \\ &= ZF_2^p + NF_2^n + n_{SRC}^A(\Delta F_2^p + \Delta F_2^n), \end{aligned}$$

where n^A_{SRC} is the number of np-SRC pairs in nucleus A, $F_2^p(x_B,Q^2)$ and $F_2^n(x_B,Q^2)$ are the free proton and neutron structure functions, $F_2^{p*}(x_B,Q^2)$ and $F_2^{n*}(x_B,Q^2)$ are the average modified structure functions for protons and neutrons in SRC pairs, and $\Delta F_2^n = F_2^{n*} - F_2^n$ (and similarly for ΔF_2^p). F_2^{p*} and F_2^{n*} are assumed to be the same for all nuclei. In this simple model, nucleon motion effects [1–3], which are also dominated by SRC pairs due to their high relative momentum, are folded into ΔF_2^p and ΔF_2^n.

This model resembles that used in [26]. However, that work focused on light nuclei and did not determine the shape of the modification function. Similar ideas using factorization were discussed in [1], such as a model-dependent ansatz for the modified structure functions which was shown to be able to describe the EMC data [27]. The analysis presented here is the first data-driven determination of the modified structure functions for nuclei from ^{3}He to lead.

Since there are no model-independent measurements of F_2^n, we apply Eq. 1 to the deuteron, rewriting F_2^n as $F_2^d - F_2^p - n^d_{SRC}\left(\Delta F_2^p + \Delta F_2^n\right)$. We then rearrange Eq. 1 to get:

$$\frac{n^d_{SRC}\left(\Delta F_2^p + \Delta F_2^n\right)}{F_2^d} = \frac{\frac{F_2^A}{F_2^d} - (Z-N)\frac{F_2^p}{F_2^d} - N}{(A/2)a_2 - N}, \qquad \text{Eq. 2}$$

where F_2^p/F_2^d was previously measured [28] and a_2 is the measured per-nucleon cross-section ratio shown by the red lines in Fig. 1b. Here we assume a_2 approximately equals the per-nucleon SRC-pair density ratio of nucleus A and deuterium: $\left(n^A_{SRC}/A\right)/\left(n^d_{SRC}/2\right)$ [1, 11-15].

Since $\Delta F_2^p + \Delta F_2^n$ is assumed to be nucleus-independent, our model predicts that the left-hand side of Eq. 2 should be a universal function (i.e., the same for all nuclei). This requires that the nucleus-dependent quantities on the right-hand side of Eq. 2 combine to give a nucleus-independent result.

This is tested in Fig. 2. The left panel shows $\left[F_2^A(x_B)/A\right] / \left[F_2^d(x_B)/2\right]$, the per-nucleon structure-function ratio of different nuclei relative to deuterium without isoscalar corrections. The approximately linear deviation from unity for $0.3 \leq x_B \leq 0.7$ is the EMC effect, which is larger for heavier nuclei. The right panel shows the relative structure modification of nucleons in *np*-SRC pairs, $n^d_{SRC}\left(\Delta F_2^p + \Delta F_2^n\right)/F_2^d$, extracted using the right-hand side of Eq. 2.

The EMC slope for all measured nuclei increases monotonically with *A* while the slope of the SRC-modified structure function is constant within uncertainties, see Fig. 3 and Extended Data Table II. Even ^{3}He, which has a dramatically different structure-function ratio due to its extreme proton-to-neutron ratio of 2, has a remarkably similar modified structure function with the same slope as the other nuclei. Thus, we conclude that the magnitude of the EMC effect in different nuclei can be described by the abundance of *np*-SRC pairs and that the proposed SRC-pair modification function is, in fact, universal. This universality appears to hold even beyond $x_B = 0.7$.

The universal function extracted here will be tested directly in the future using lattice QCD calculations [26] and by measuring semi-inclusive DIS off the deuteron, tagged by the detection of a high-momentum backward-recoiling proton or neutron that will allow to directly quantify the relationship between the momentum and the structure-function modification of bound nucleons [29].

The universal SRC-pair modification function can also be used to extract the free neutron-to-proton structure-function ratio, F_2^n/F_2^p, by applying Eq. 1 to the deuteron and using the measured proton and deuteron structure functions (see Extended Data Fig. 1). In addition to its own importance, this F_2^n can be used to apply self-consistent isoscalar corrections to the EMC effect data (see Supplementary Information Eq. 5).

To further test the SRC-driven EMC model, we consider the isophobic nature of SRC pairs (i.e., *np*-dominance), which leads to an approximately constant probability for a neutron to belong to an SRC pair in medium to heavy nuclei, while the proton probability increases as *N*/*Z* [22]. If the EMC effect is indeed driven by high-momentum SRCs, then in neutron-rich nuclei both the neutron EMC effect and the SRC probability should saturate, while for protons both should grow with the nuclear mass and the neutron excess.

This is done by examining the correlation of the individual per-proton and per-neutron QE SRC cross-section ratios, $a_2^p = (\sigma_A/Z)/\sigma_d$ and $a_2^n = (\sigma_A/N)/\sigma_d$, and DIS EMC slopes, dR^p_{EMC}/dx_B and dR^n_{EMC}/dx_B (see Extended Data Tables I and III and Supplementary Information sections III and V).

Figure 4 shows the per-proton and per-neutron EMC slopes as a function of a_2^p and a_2^n, respectively. We consider these correlations both before (top panels) and after (bottom panels) applying isoscalar corrections to the EMC data and compare them with the predictions of the SRC-driven EMC model. By not applying isoscalar corrections, the top panel allows focusing on the separate behavior of protons and neutrons. Applying self-consistent isoscalar corrections makes both the per-neutron and per-proton EMC-SRC correlations linear, in overall agreement with the model prediction for $N = Z$ nuclei.

This simple rescaling of the previous EMC-SRC correlation result [4, 5], as expected, does not change the EMC-SRC correlation or its slope. However, the per-neutron and per-proton results differ significantly. Because the probability that a neutron belongs to an SRC pair does not increase for nuclei heavier than C ($A = 12$) [22], our model predicts that the per-neutron EMC effect (i.e., the slope of $\frac{F_2^A/N}{F_2^d/1}$) will also not increase for $A \geq 12$. In contrast, the probability that a proton belongs to an SRC pair continues to increase for all measured nuclei [22] and therefore the per-proton EMC effect should continue to increase for all measured nuclei. This saturation / no-saturation is a non-trivial prediction of our model that is supported by the data.

In the per-neutron correlation, the proton-rich ^{3}He point is far below the simple straight line, while the neutron-rich Fe and Pb points are above it. In the per-proton correlation, the proton-rich ^{3}He point is below the simple straight line for $N = Z$ nuclei, while the increasingly neutron-rich heavy nuclei are above it. These features of the data are all well-described by our SRC-driven EMC model.

To conclude, the association of the EMC effect with SRC pairs implies that it is a dynamical effect. Most of the time, nucleons bound in nuclei have the same internal structure as that of free nucleons. However, for short time intervals when two nucleons form a temporary high local-density SRC pair, their internal structure is briefly modified. When the two nucleons disassociate, their internal structure again becomes similar to that of free nucleons. This dynamical picture differs significantly from the traditional static modification in the nuclear mean-field, previously proposed as an explanation for the EMC effect.

The new universal modification function presented here has implications for our understanding of fundamental aspects of Quantum Chromodynamics (QCD). For example, the study of the ratio of the d-quark to u-quark population in a free nucleon as $x_B \rightarrow 1$ offers a stringent test of symmetry-breaking mechanisms in QCD. This can be extracted from measuring the free proton to neutron structure-function ratio. However, the lack of a free neutron target forces the use of proton and deuterium DIS data, which requires corrections for the deuteron EMC effect to extract the free neutron. The universal SRC modification function presented here does just that, in a data-driven manner, see Extended Data Fig. 1.

Turning to neutron-rich nuclei, the larger proton EMC effect has several implications. As the proton has two u-quarks and one d-quark while the neutron has two d-quarks and one u-quark, the larger average modification of the protons' structure implies a larger average modification of the distribution of u-quarks in the nucleus as compared to d-quarks. This will affect DIS charge-changing neutrino interactions, because neutrinos (ν) scatter preferentially from d-quarks and anti-neutrinos ($\bar{\nu}$) from u-quarks. Different modifications to d and u quark distributions will cause a difference in the ν and $\bar{\nu}$ cross-sections in asymmetric nuclei, which could then be misinterpreted as a sign of physics beyond the standard model or of CP-violation. One example of this is the NuTeV experiment, which extracted an anomalous value of the standard-model Weinberg mixing angle from ν and $\bar{\nu}$-nucleus DIS on iron. Ref. [30] pointed out that this anomaly could be due to differences between the proton and the neutron caused by mean-field effects. Our model provides an alternative mechanism. Similarly, the future DUNE experiment will use high-energy ν and $\bar{\nu}$ beams incident on the asymmetric nucleus ^{40}Ar to look for differences in ν and $\bar{\nu}$ oscillations as a possible mechanism for explaining the matter-antimatter asymmetry. They will therefore also need to take the larger proton EMC effect into account to avoid similar anomalies.

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Figure Captions

Fig 1 | DIS and QE (e,e′) Cross-section Ratios. The per-nucleon cross-section ratios of nucleus with atomic number A to deuterium for (a. 1 - 4) DIS kinematics ($0.2 \leq x_B \leq 0.6$ and $W \geq 1.8$ GeV). The solid points show the data of this work, the open squares the data of [9] and the open triangles show the data of [10]. The red lines show the linear fit. (b. 1 - 4) QE kinematics ($0.8 \leq x_B \leq 1.9$). The solid points show the data of this work and the open squares the data of [11]. The red lines show the constant fit. The error bars shown include both statistical and point-to-point systematic uncertainties, both at the 1σ or 68% confidence level. The data are not isoscalar corrected.

Fig 2 | Universality of SRC pair quark distributions. The EMC effect for different nuclei, as observed in (a) ratios of $(F_2^A/A)/(F_2^d/2)$ as a function of x_B and (b) the modification of SRC pairs, as described by Eq. 2. Different colors correspond to different nuclei, as indicated by the color scale on the right. The open circles show SLAC data [9] and the open squares show Jefferson Lab data [10]. The nucleus-independent (universal) behavior of the SRC modification, as predicted by the SRC-driven EMC model, is clearly observed. The error bars on the symbols show both statistical and point-to-point systematic uncertainties, both at the 1σ or 68% confidence level and the gray bands show the median normalization uncertainty. The data are not isoscalar corrected.

Fig 3 | EMC and universal modification function slopes. The slopes of the EMC effect for different nuclei from Fig. 2a (blue) and of the universal function from Fig. 2b (red). The error bars shown include the fit uncertainties at the 1σ or 68% confidence level.

Fig 4 | Growth and saturation of the EMC effect for protons and neutrons. The (a) per-neutron and (b) per-proton strength of the EMC effect versus the corresponding per-neutron and per-proton number of SRC pairs. New data are shown by squares and existing data by circles. The dashed line shows the results of Eq. 2 using the universal modification function shown in Fig. 2 for symmetric N = Z nuclei. The solid line shows the same results for the actual nuclei. The gray region shows the effects of per-neutron saturation. (c) and (d): the same, but with isoscalar corrections. The error bars on the symbols show both statistical and systematic uncertainties, both at the 1σ or 68% confidence level.

Methods

Experimental setup and electron identification. CLAS used a toroidal magnetic field with six sectors of drift chambers, scintillation counters, Cerenkov counters and electromagnetic calorimeters to identify electrons and reconstruct their trajectories [23].

The experiment used a specially designed double target setup, consisting of a 2-cm long cryo-target cell, containing liquid deuterium, and a solid target [24]. The cryo-target cell and solid target were separated by 4 cm, with a thin isolation foil between them. Both targets and the isolation foil were kept in the beam line simultaneously. This allowed for an accurate measurement of cross-section ratios for nuclei relative to deuterium. A dedicated control system was used to position one of six different solid targets (thin and thick Al, Sn, C, Fe, and Pb, all in natural abundance) at a time during the experiment. The main data collected during the experiment was for a target configuration of deuterium + C, Fe, or Pb and also for an empty cryo-target cell with the thick Al target.

We identified electrons by requiring that the track originated in the liquid deuterium or solid targets, produced a large enough signal in the Cerenkov counter, and deposited enough energy in the Electromagnetic Calorimeter, see [21, 22] for details.

Vertex reconstruction. Electrons scattering from the solid and cryo-targets were selected using vertex cuts with a resolution of several mm (depending on the scattering angle), which is sufficient to separate the targets which are 4 cm apart [21]. We considered events with reconstructed electron vertex up to 0.5 cm outside the 2 cm long cryo-target to originate from the deuterium. Similarly, for the solid target, we considered events with reconstructed electron vertex up to 1.5 cm around it.

Background subtraction. There are two main sources of background in the measurement: (1) electrons scattering from the Al walls of the cryo-target cell, (2) electrons scattering from the isolation foil between the cryo-target and solid target. When the vertex of these electrons is reconstructed within the region of the deuterium target, they falsely contribute to the cross section associated with the deuterium target. Data from measurements done using an empty cryo-target is used to subtract these contributions. In the case of QE scattering, at $x_B > 1$, these measurements do not have enough statistics to allow for a reliable background subtraction. We therefore require QE deuterium electrons to be reconstructed in the inner 1-cm of the 2-cm long cryo-target. This increases the reliability of the background subtraction but reduces the deuterium statistics by a factor of two.

Data from runs with a full cryo-target and no solid target were used to subtract background from electron scattering events with a reconstructed vertex in the solid-target region, originating from the isolation foil or the cryo-target.

To increase statistics, the analysis combined all deuterium data, regardless of the solid target placed with it in the beam line. We only consider runs where the electron scattering rate from the cryo-target deviated by less than 4% from the average.

The systematic uncertainties associated with the vertex cuts, target wall subtraction, and combination of deuterium data from different runs are described in the Supplemental Materials, section 2.

Data Availability: The raw data from this experiment are archived in Jefferson Lab's mass storage silo.

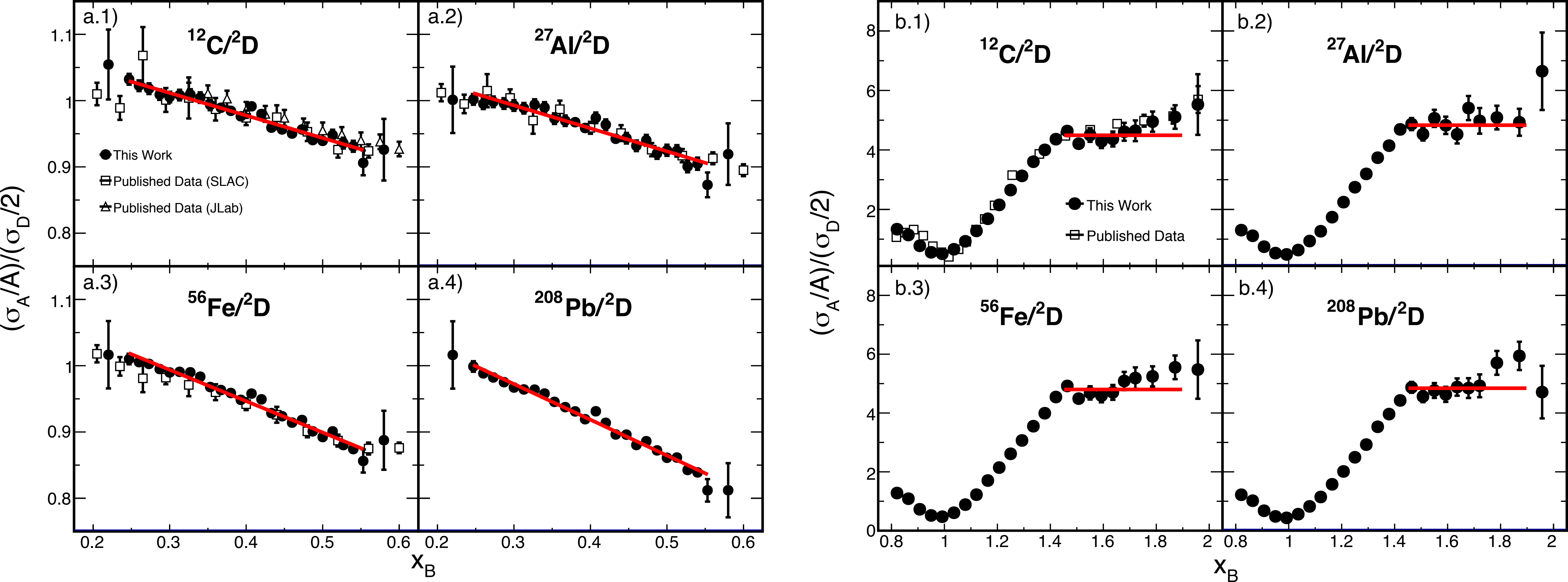
Extended Data Figure and Tables Captions

Extended Data Fig 1 | F_2^n/F_2^p Models. The ratio of neutron to proton structure functions, F_2^n/F_2^p, derived from the SRC-driven EMC model (blue band), assumed in the isoscalar corrections of Refs. [9] (red line) and [10] (green line), and derived in the CT14 global fit, shown here for Q^2 = 10 GeV^2 (gray band). The large spread among the various models shows the uncertainty in F_2^n, a key ingredient in the isoscalar corrections previously applied to the EMC effect data

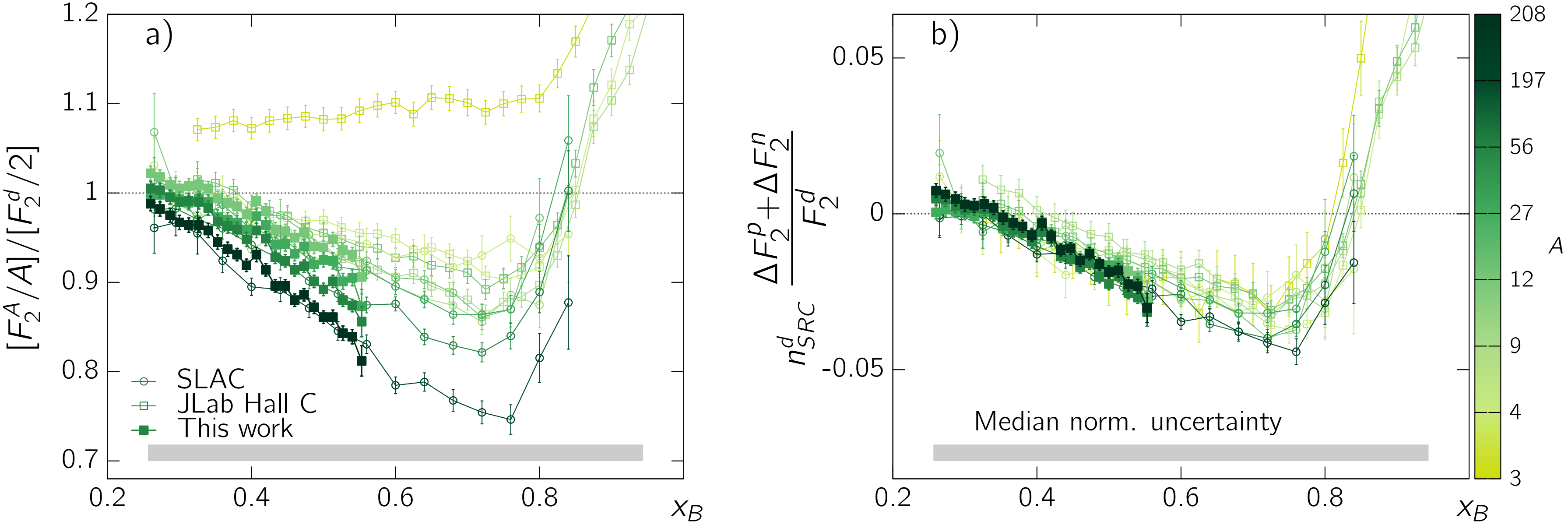
Extended Data Table I: | SRC Scaling Coefficients. Per-nucleon (a_2), per-proton (a_2^p), and per-neutron (a_2^n) SRC scale factors for nucleus *A* relative to deuterium. The 1σ or 68% confidence level uncertainties shown include the fit uncertainties.

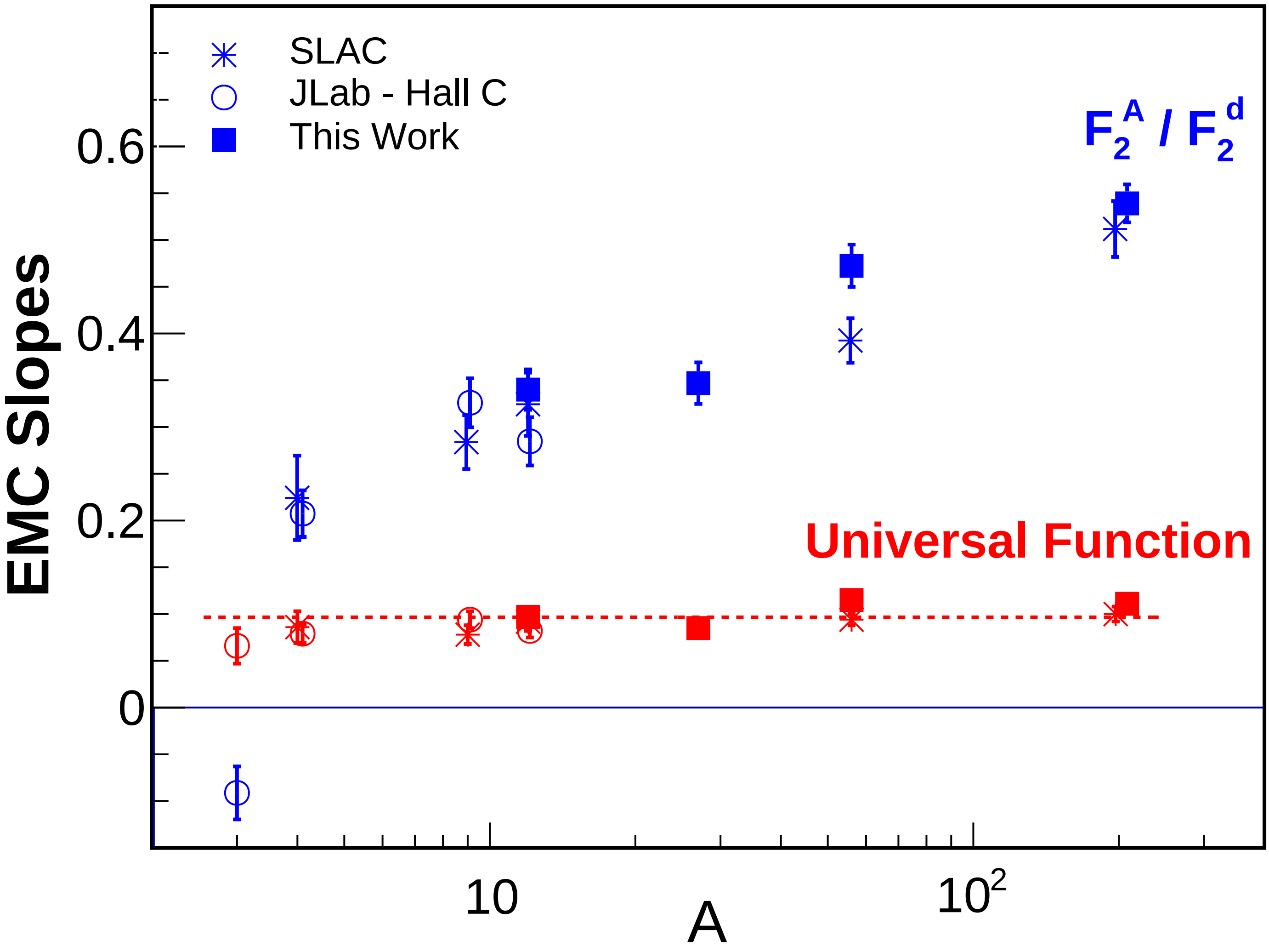
Extended Data Table II: | EMC Slopes. Slopes of non isoscalar-corrected F_2^A/F_2^d (dR_{EMC}/dx_B) and the universal function, shown in Figs. 2a and 2b of the main paper, respectively. The SLAC data is from [9] and the JLab Hall C data is from [10]. The slopes are obtained from a linear fit of the data for $0.25 \leq x_B \leq 0.7$. The 1σ or 68% confidence level uncertainties shown include the fit uncertainties.

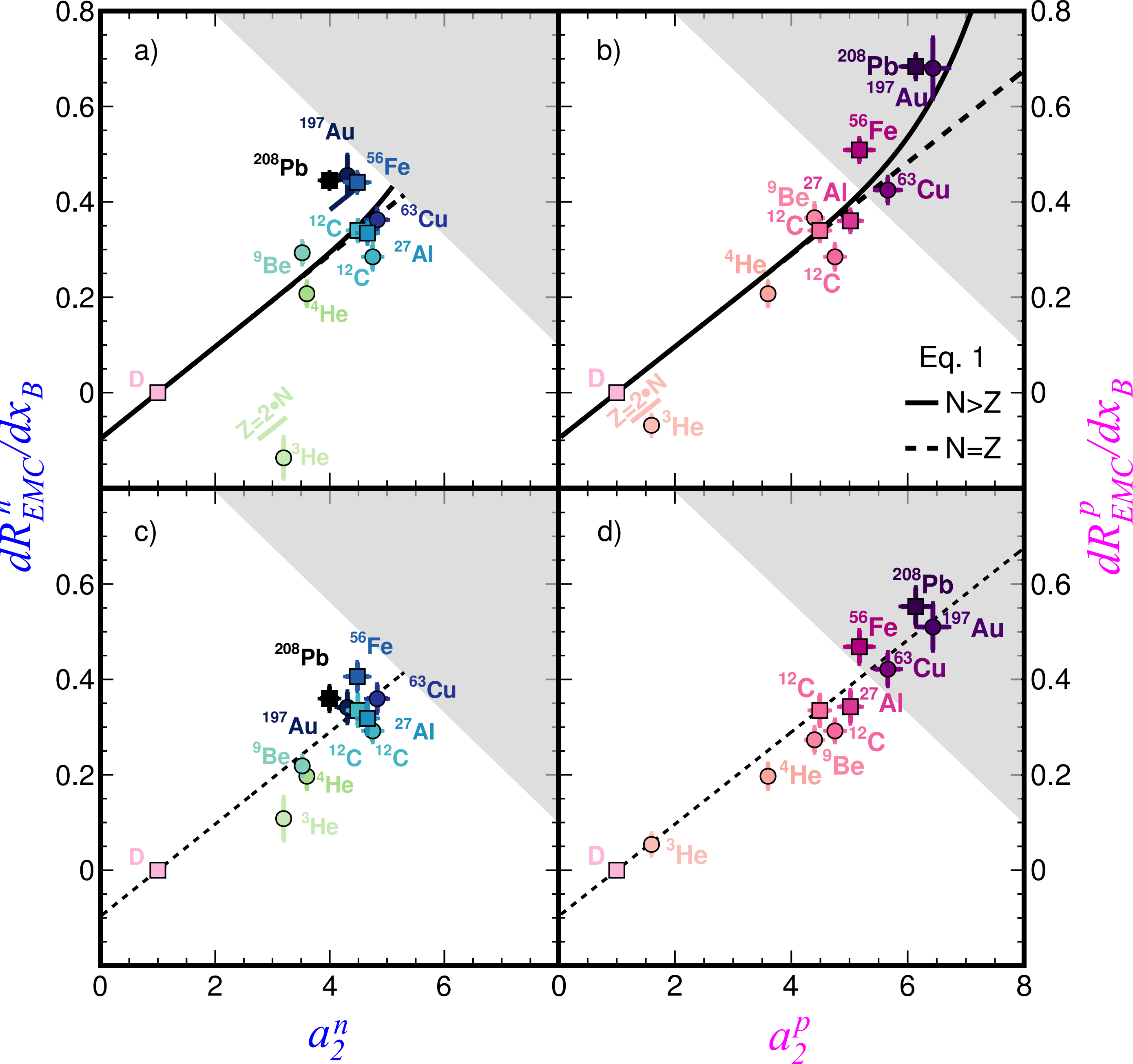
Extended Data Table III: | Per nucleon, per-proton, and per-neutron EMC Slopes. Per-nucleon (dR_{EMC}/dx_B) per-proton (dR_{EMC}^p/dx_B) and per-neutron (dR_{EMC}^n/dx_B) EMC slopes from the current and previous works, used in Fig. 4 of the main paper. The previous data shows the JLab Hall C results [10] for light nuclei ($A \leq 12$) and the SLAC results [9] for heavier nuclei. The 1σ or 68% confidence level uncertainties shown include the fit uncertainties.



a.1)
^{12}C/^{2}D
a.2)
^{27}Al/^{2}D
a.3)
^{56}Fe/^{2}D
a.4)
^{208}Pb/^{2}D
● This Work
□ Published Data (SLAC)
△ Published Data (JLab)
1.1
1
0.9
0.8
0.2
0.3
0.4
0.5
0.6
$(\sigma_A/A)/(\sigma_D/2)$
x_B
b.1)
^{12}C/^{2}D
b.2)
^{27}Al/^{2}D
b.3)
^{56}Fe/^{2}D
b.4)
^{208}Pb/^{2}D
● This Work
□ Published Data
8
6
4
2
0
0.8
1
1.2
1.4
1.6
1.8
2
$(\sigma_A/A)/(\sigma_D/2)$
x_B


a)
$[F_2^A/A]/[F_2^d/2]$
1.2
1.1
1
0.9
0.8
0.7
0.2
0.4
0.6
0.8
x_B
SLAC
JLab Hall C
This work
b)
$\frac{\Delta F_2^p + \Delta F_2^n}{n_{SRC}^d F_2^d}$
0.05
0
-0.05
0.2
0.4
0.6
0.8
x_B
Median norm. uncertainty
208
197
56
27
12
9
4
3
A


SLAC
JLab - Hall C
This Work
F_2^A / F_2^d
Universal Function
EMC Slopes
0.6
0.4
0.2
0
10
10^2
A


a)
b)
c)
d)
dR^n_{EMC}/dx_B
dR^p_{EMC}/dx_B
a^n_2
a^p_2
0.8
0.6
0.4
0.2
0
0
2
4
6
8
^{208}Pb
^{197}Au
^{56}Fe
^{63}Cu
^{27}Al
^{12}C
^{9}Be
^{4}He
^{3}He
D
Z=2•N
Eq. 1
N>Z
N=Z